

Differential operators on supercircle: conformally equivariant quantization and symbol calculus

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Abstract

We consider the supercircle $S^{1|1}$ equipped with the standard contact structure. The conformal Lie superalgebra $\mathcal{K}(1)$ acts on $S^{1|1}$ as the Lie superalgebra of contact vector fields; it contains the Möbius superalgebra $\mathfrak{osp}(1|2)$. We study the space of linear differential operators on weighted densities as a module over $\mathfrak{osp}(1|2)$. We introduce the canonical isomorphism between this space and the corresponding space of symbols and find interesting resonant cases where such an isomorphism does not exist.

1 Introduction

Conformally and projectively equivariant symbol calculus and quantization (see, e.g., [18, 6]) is the canonical way to define a “total symbol” of a differential operator on a manifold equipped with a G -structure (e.g., projective or conformal structure). In the case of supermanifolds, the only information available at the moment is for the supercircle $S^{1|1}$ with a contact structure, see [4].

The Lie superalgebra, $\mathcal{K}(1)$ of contact vector fields on $S^{1|1}$ is often called a *superconformal algebra*, cf. [14]. This algebra and its central extension are the simplest super-generalizations of the Witt and Virasoro algebras, respectively. The Lie superalgebra $\mathfrak{osp}(1|2)$ plays on $S^{1|1}$ the same crucial role that $\mathfrak{sl}(2)$ plays on S^1 .

The main object of our study is the space of linear differential operators acting on weighted densities. The space \mathcal{F}_λ of weighted densities with weight $\lambda \in \mathbb{C}$ (or λ -densities for short) is a module over $\mathcal{K}(1)$. Therefore, the space $\mathcal{D}_{\lambda,\mu}$ of linear differential operators from \mathcal{F}_λ to \mathcal{F}_μ is also a $\mathcal{K}(1)$ -module. Each module $\mathcal{D}_{\lambda,\mu}$ has a natural filtration by the order of differential operators; the graded module $\text{gr}\mathcal{D}_{\lambda,\mu}$ is called the *space of symbols*. We restrict the $\mathcal{K}(1)$ -module structures to a particular subalgebra $\mathfrak{osp}(1|2)$ and look for $\mathfrak{osp}(1|2)$ -isomorphisms

$$\sigma : \mathcal{D}_{\lambda,\mu} \rightarrow \text{gr}\mathcal{D}_{\lambda,\mu}, \quad Q : \text{gr}\mathcal{D}_{\lambda,\mu} \rightarrow \mathcal{D}_{\lambda,\mu},$$

where $Q = \sigma^{-1}$. These isomorphisms are called the $\mathfrak{osp}(1|2)$ -equivariant *symbol map* and *quantization map*, respectively.

For almost all values (λ, μ) , we prove the existence and uniqueness (up to normalization) of the $\mathfrak{osp}(1|2)$ -equivariant symbol map σ and calculate its explicit formula. We also calculate its inverse $Q = \sigma^{-1}$ generalizing the quantization map from [4].

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We consider particular values of the weights that we call *resonant* which satisfy

$$\mu - \lambda = \frac{1}{2}, 1, \frac{3}{2}, 2, \dots \quad (1.1)$$

We prove that for these values there is no $\mathfrak{osp}(1|2)$ -isomorphism between the space of linear differential operators and the corresponding space of symbols, except for the special values:

$$\lambda = \frac{1-m}{4}, \quad \mu = \frac{1+m}{4}, \quad (1.2)$$

where m is odd. We calculate the cohomological obstructions to existence of such an isomorphism.

The $\mathcal{K}(1)$ -modules $\mathcal{D}_{\lambda,\mu}$, where (λ, μ) as in (1.2), are of particular interest. These modules are characterized by the existence of $\mathfrak{osp}(1|2)$ -invariants. Let us mention that in the classical case of S^1 the corresponding modules are closely related to the Virasoro and Adler-Gelfand-Dickey algebras, see, e.g., [24].

Our main tool is the *finer filtration* on the space of differential operators $\mathcal{D}_{\lambda,\mu}$:

$$\mathcal{D}_{\lambda,\mu}^0 \subset \mathcal{D}_{\lambda,\mu}^{\frac{1}{2}} \subset \mathcal{D}_{\lambda,\mu}^1 \subset \mathcal{D}_{\lambda,\mu}^{\frac{3}{2}} \subset \dots \subset \mathcal{D}_{\lambda,\mu}^{\ell-\frac{1}{2}} \subset \mathcal{D}_{\lambda,\mu}^\ell \subset \dots \quad (1.3)$$

We introduce the notion of differential operators of *semi-integer order*. In particular, the space $\text{Vect}_{\mathbb{C}}(S^{1|1})$ of all vector fields (2.2) is, as $\mathcal{K}(1)$ -module, a direct sum of two submodules: $\mathcal{K}(1)$ itself and the space of tangent vector fields. The finer filtration is stable with respect to the $\mathcal{K}(1)$ -action.

It worth noticing that the results of this paper remain true in the case of pseudodifferential operators, as considered in [4], but we will not dwell on it.

2 Geometry of the supercircle

The *supercircle* $S^{1|1}$ is the simplest supermanifold of dimension $1|1$ generalizing S^1 . In order to fix notation, let us give here the basic definitions of geometric objects on $S^{1|1}$; for more details, see [2, 19, 21, 15].

We define the supercircle $S^{1|1}$ in terms of its superalgebra of functions, denoted by $C_{\mathbb{C}}^\infty(S^{1|1})$ and consisting of elements of the form

$$f(x, \xi) = f_0(x) + \xi f_1(x) \quad (2.1)$$

where x is an arbitrary parameter on S^1 and ξ is an (formal Grassmann) coordinate such that $\xi^2 = 0$. The parity function p is defined by $p(f(x)) = 0$ and $p(\xi) = 1$.

2.1 Vector fields and differential forms

A *vector field* on $S^{1|1}$ is a superderivation of $C_{\mathbb{C}}^\infty(S^{1|1})$. Every vector field can be expressed in coordinates in terms of partial derivatives:

$$X = f \frac{\partial}{\partial x} + g \frac{\partial}{\partial \xi}, \quad (2.2)$$

where $f, g \in C_{\mathbb{C}}^\infty(S^{1|1})$. The space of vector fields is a Lie superalgebra denoted by $\text{Vect}_{\mathbb{C}}(S^{1|1})$. Let $\Omega^1(S^{1|1})$ be the rank $1|1$ right $C_{\mathbb{C}}^\infty(S^{1|1})$ -module with basis dx and $d\xi$; we interpret it as the *right* dual over $C_{\mathbb{C}}^\infty(S^{1|1})$ to the left $C_{\mathbb{C}}^\infty(S^{1|1})$ -module $\text{Vect}_{\mathbb{C}}(S^{1|1})$, by setting $\langle \partial_{y_i}, dy_j \rangle = \delta_{ij}$ for $y = (x, \xi)$. The space $\Omega^1(S^{1|1})$ is a *left* module over $\text{Vect}_{\mathbb{C}}(S^{1|1})$, the action being given by the *Lie derivative*:

$$\langle X, L_Y \alpha \rangle := \langle [X, Y], \alpha \rangle.$$

2.2 Lie superalgebra of contact vector fields

The standard *contact structure*¹ on $S^{1|1}$ is defined as a codimension 1 *non-integrable* distribution $\langle \overline{D} \rangle$ on $S^{1|1}$, i.e., a subbundle in $TS^{1|1}$ generated by the odd vector field

$$\overline{D} = \frac{\partial}{\partial \xi} - \xi \frac{\partial}{\partial x}. \quad (2.3)$$

This contact structure can be equivalently defined as the kernel of the differential 1-form

$$\alpha = dx + \xi d\xi. \quad (2.4)$$

A vector field X on $S^{1|1}$ is said to be *contact* if it preserves the contact distribution:

$$[X, \overline{D}] = \psi_X \overline{D},$$

where $\psi_X \in C_{\mathbb{C}}^{\infty}(S^{1|1})$ is a function depending on X . The space of contact vector fields is a Lie superalgebra denoted by $\mathcal{K}(1)$. The following statement is well-known.

Lemma 2.1. *Every contact vector field can be expressed, for any $f \in C_{\mathbb{C}}^{\infty}(S^{1|1})$, as²*

$$X_f = -f \overline{D}^2 + \frac{1}{2} D(f) \overline{D}, \quad \text{where} \quad D = \frac{\partial}{\partial \xi} + \xi \frac{\partial}{\partial x}. \quad (2.5)$$

The vector field (2.5) is said to be the contact vector field with *contact Hamiltonian* f . One checks that

$$L_{X_f} \alpha = f' \alpha, \quad [X_f, \overline{D}] = -\frac{1}{2} f' \overline{D}. \quad (2.6)$$

The *contact bracket* is defined by $[X_f, X_g] = X_{\{f, g\}}$. The space $C_{\mathbb{C}}^{\infty}(S^{1|1})$ is thus equipped with a Lie superalgebra structure isomorphic to $\mathcal{K}(1)$. The explicit formula can be easily calculated:

$$\{f, g\} = fg' - f'g + (-1)^{p(f)(p(g)+1)} \frac{1}{2} D(f)D(g). \quad (2.7)$$

2.3 Projective/conformal symmetries: $\mathfrak{osp}(1|2)$ -action

In the case of S^1 , the notions of *projective* and *conformal* structures coincide and are defined by the action of $\mathfrak{sl}(2)$. In the adapted (local) coordinate x on S^1 this action is spanned by three vector fields:

$$\mathfrak{sl}(2) = \text{Span} \left(\frac{\partial}{\partial x}, \quad x \frac{\partial}{\partial x}, \quad x^2 \frac{\partial}{\partial x} \right) \quad (2.8)$$

corresponding to the fraction-linear transformations

$$x \mapsto \frac{ax + b}{cx + d}, \quad ad - bc = 1.$$

A *projective structure* on S^1 is given by an atlas with fraction-linear coordinate transformations (in other words, by an atlas such that the $\mathfrak{sl}(2)$ -action (2.8) is well-defined). Classification of projective structures on S^1 is equivalent to classification of coadjoint orbits of the Virasoro algebra (see [16], [24]).

¹This structure is famous in mathematical physics, it is also known as the “SUSY-structure”.

²For interpretation of the fields D and \overline{D} , see [26].

A projective/conformal structure on $S^{1|1}$ is defined as a (local) action of the orthosymplectic Lie superalgebra $\text{osp}(1|2)$ generated by D and xD :

$$\begin{aligned} \text{osp}(1|2)_0 &= \text{Span} \left(X_1 = \frac{\partial}{\partial x}, \quad X_x = x \frac{\partial}{\partial x} + \frac{1}{2} \xi \frac{\partial}{\partial \xi}, \quad X_{x^2} = x^2 \frac{\partial}{\partial x} + x \xi \frac{\partial}{\partial \xi} \right), \\ \text{osp}(1|2)_1 &= \text{Span} (X_\xi = D, \quad X_{\xi x} = x D). \end{aligned} \quad (2.9)$$

Remark 2.2. The $\text{osp}(1|2)$ -action (2.9) is the infinitesimal version of the contact fractional-linear transformations

$$(x, \xi) \mapsto \left(\frac{ax + b + \gamma \xi}{cx + d + \delta \xi}, \frac{\alpha x + \beta + e \xi}{cx + d + \delta \xi} \right), \quad (2.10)$$

where $ad - bc - \alpha\beta = 1$, $e^2 + 2\gamma\delta = 1$, $\alpha e = a\delta - c\gamma$ and $\beta e = b\delta - d\gamma$ (cf. [4]).

As in the S^1 case, these contact fraction-linear transformations preserve the action (2.9) so that atlases compatible with the projective structure are precisely those atlases for which the action (2.9) is well defined. For a classification of projective/conformal structures on $S^{1|1}$, also equivalent to classification of the orbits in the coadjoint representations of the Neveu-Schwarz and Ramond superalgebras, see [23].

Remark 2.3. Eq. (2.6) implies that the (super)centralizer of \overline{D} is spanned by $X_\xi = D$ and $X_1 = D^2$.

3 Modules of weighted densities

We introduce a 1-parameter family of modules over the Lie superalgebra $\mathcal{K}(1)$. As vector spaces all these modules are isomorphic to $C_{\mathbb{C}}^\infty(S^{1|1})$, but not as $\mathcal{K}(1)$ -modules.

For every contact vector field X_f , define a 1-parameter family of first-order differential operators on $C_{\mathbb{C}}^\infty(S^{1|1})$:

$$L_{X_f}^\lambda = X_f + \lambda f', \quad \lambda \in \mathbb{C}. \quad (3.1)$$

One easily checks the the map $X_f \mapsto L_{X_f}^\lambda$ is a homomorphism of Lie superalgebras, i.e., $[L_{X_f}^\lambda, L_{X_g}^\lambda] = L_{[X_f, X_g]}^\lambda$, for every λ . One thus obtains a 1-parameter family of $\mathcal{K}(1)$ -modules on $C_{\mathbb{C}}^\infty(S^{1|1})$ that we denote \mathcal{F}_λ and call the space of *densities of weight λ* (or λ -densities for short). The space of 1-forms proportional to α is the space of sections of the line bundle $\langle \overline{D} \rangle^\perp \subset T_{\mathbb{C}}^* S^{1|1}$ over $S^{1|1}$, i.e., the line bundle of covectors orthogonal to the contact distribution. A λ -density is a section of the line bundle $\left(\langle \overline{D} \rangle^\perp \right)^{\otimes \lambda}$. It is then natural to express every λ -density in terms of the contact form α as

$$\phi = f \alpha^\lambda, \quad \text{where } f \in C^\infty(S^{1|1}).$$

Example 3.1. (a) The module \mathcal{F}_0 is nothing but the space of functions $C^\infty(S^{1|1})$.

(b) The module \mathcal{F}_1 is the space of 1-forms proportional to α .

(c) A more interesting example is provided by the Lie superalgebra $\mathcal{K}(1) \cong \mathcal{F}_{-1}$ viewed as a module over itself, see below.

As $\mathcal{K}(1)$ -module, the space of volume forms on $S^{1|1}$ is isomorphic (up to parity, perhaps) to $\mathcal{F}_{\frac{1}{2}}$. Therefore, *Berezin integral* ([2, 19]) $\mathcal{B} : \mathcal{F}_{\frac{1}{2}} \rightarrow \mathbb{C}$ can be given, for any $f = f_0(x) + \xi f_1(x)$, by the formula

$$\mathcal{B}(f \alpha^{\frac{1}{2}}) := \int_{S^1} f_1(x) dx.$$

So the product of densities composed with \mathcal{B} yields a bilinear $\mathcal{K}(1)$ -invariant form:

$$\langle \cdot, \cdot \rangle : \mathcal{F}_\lambda \otimes \mathcal{F}_\mu \rightarrow \mathbb{C}, \quad \lambda + \mu = \frac{1}{2}. \quad (3.2)$$

It turns out that the adjoint $\mathcal{K}(1)$ -module is isomorphic to \mathcal{F}_{-1} . In other words, a contact Hamiltonian is a density of weight -1 rather than a function:

$$\{f, g\} = L_{X_f}^{-1} g. \quad (3.3)$$

This statement can be reformulated as follows.

Corollary 3.2. *For a contact vector field X_f given by (2.5), the expression*

$$\sigma(X_f) := f \alpha^{-1} \quad (3.4)$$

is a well defined -1 -density. The expression (3.4) is $\mathcal{K}(1)$ -invariant and independent of the choice of the contact form α .

3.1 Poisson algebra of weighted densities

The contact bracket (2.7) extends to densities of arbitrary weight: $\{, \} : \mathcal{F}_\lambda \otimes \mathcal{F}_\mu \rightarrow \mathcal{F}_{\lambda+\mu+1}$ and defines a structure of Poisson Lie superalgebra. Explicitly ([20])

$$\{f, g\} = \lambda f g' - \mu f' g + (-1)^{p(f)(p(g)+1)} \frac{1}{2} D(f) D(g). \quad (3.5)$$

The following statement can be checked directly (cf. Grozman's list of invariant operators [13]).

Proposition 3.3. *The operation (3.5) is $\mathcal{K}(1)$ -invariant and satisfies the Jacobi and Leibniz identities makes the space of weighted densities on $S^{1|1}$ a Poisson superalgebra.*

3.2 Splitting of vector fields

Proposition 3.4. *There is an isomorphism of $\mathcal{K}(1)$ -modules*

$$\text{Vect}(S^{1|1}) \cong \mathcal{F}_{-\frac{1}{2}} \oplus \mathcal{F}_{-1}. \quad (3.6)$$

Proof. The submodule $\mathcal{F}_{-1} \subset \text{Vect}_{\mathbb{C}}(S^{1|1})$ is the subalgebra $\mathcal{K}(1)$ itself (see Corollary 3.2); the submodule $\mathcal{F}_{-\frac{1}{2}} \subset \text{Vect}_{\mathbb{C}}(S^{1|1})$, consists of the vector fields tangent to the contact distribution: $X = g \overline{D}$, where g is an arbitrary function. \square

Note that an analog of Proposition 3.4 holds for any contact (super)manifold (see [22]).

4 Differential operators on weighted densities

For differential operators on supermanifolds, see [3, 21, 4, 1].

4.1 The modules $\mathcal{D}_{\lambda,\mu}$: definition

The classical Peetre theorem [25] implies that the following two definitions are equivalent. They reflect two different aspects of the notion: geometric and algebraic.

1) A linear map $A : \mathcal{F}_\lambda \rightarrow \mathcal{F}_\mu$ is called a *differential operator* if it is *local*, that is, if it preserves the supports of the arguments: $\text{Supp} A(\phi) \subset \text{Supp} \phi$, where the support of a λ -density is a subset of S^1 given by the union of the supports of the even and the odd parts.

2) A linear map $A : \mathcal{F}_\lambda \rightarrow \mathcal{F}_\mu$ is a *differential operator* if there exists an integer k such that $[\cdots [A, f_1], \cdots, f_{k+1}] = 0$ for any functions f_1, \dots, f_{k+1} . The minimal such k is called the *order* of A .

We denote $\mathcal{D}_{\lambda,\mu}$ the space of linear differential operators from \mathcal{F}_λ to \mathcal{F}_μ and $\mathcal{D}_{\lambda,\mu}^k$ the space of linear differential operators of order k . One has a filtration

$$\mathcal{D}_{\lambda,\mu}^0 \subset \mathcal{D}_{\lambda,\mu}^1 \subset \cdots \subset \mathcal{D}_{\lambda,\mu}^k \subset \cdots$$

Example 4.1. The space of zeroth-order operators is $\mathcal{D}_{\lambda,\mu}^0 \cong \mathcal{F}_{\mu-\lambda}$; it consists of the operators of multiplication by $(\mu - \lambda)$ -densities. The space of first-order operators on densities of fixed weight $\lambda = \mu$ is $\mathcal{D}_{\lambda,\lambda}^1 \cong \text{Vect}_{\mathbb{C}}(S^{1|1}) \oplus C_{\mathbb{C}}^\infty(S^{1|1})$.

The space $\mathcal{D}_{\lambda,\mu}$ is naturally a module over $\mathcal{K}(1)$; the action being given by the commutator with Lie derivative:

$$\mathcal{L}_{X_f}^{\lambda,\mu}(A) := L_{X_f}^\mu \circ A - A \circ L_{X_f}^\lambda. \quad (4.1)$$

The above filtration is obviously $\mathcal{K}(1)$ -invariant.

Proposition 4.2. (i) Every differential operator $A \in \mathcal{D}_{\lambda,\mu}$ can be expressed in the form:

$$A(f \alpha^\lambda) = \sum_{i=0}^{\ell} a_i(x, \xi) \overline{D}^i(f) \alpha^\mu, \quad (4.2)$$

where the coefficients $a_i(x, \xi)$ are arbitrary functions and $\ell \in \mathbb{N}$.

(ii) If $A \in \mathcal{D}_{\lambda,\mu}^k$, then $\ell = 2k$.

Proof. Part (i). Every DO on S^1 is a (finite) expression $A = \sum_{i \geq 0} a_i(x) \left(\frac{\partial}{\partial x} \right)^i$, see [25]. It follows

that a differential operator on $S^{1|1}$ can be expressed as a 2×2 -matrix of differential operators on S^1 , or, equivalently,

$$A = \sum_{i \geq 0} \tilde{a}_i(x, \xi) \left(\frac{\partial}{\partial x} \right)^i + \sum_{i \geq 0} \tilde{b}_i(x, \xi) \left(\frac{\partial}{\partial x} \right)^i \frac{\partial}{\partial \xi}.$$

Since $\frac{\partial}{\partial x} = -\overline{D}^2$, every differential operator A is a polynomial expression in \overline{D} .

Part (ii) is straightforward. □

Remark 4.3. (a) There is another way to express a differential operator in local coordinates:

$$A(f \alpha^\lambda) = \sum_{i \geq 0} b_i(x, \xi) D^i(f) \alpha^\mu. \quad (4.3)$$

where the coefficients b_i are related to a_i via

$$b_{2i} = (-1)^i (a_{2i} - 2\xi a_{2i-1}), \quad b_{2i+1} = (-1)^i a_{2i+1}. \quad (4.4)$$

Formula (4.4) follows from the expression $\overline{D} = D - 2\xi D^2$.

(b) Clearly, $b(x, \xi) D^i \in \mathcal{D}_{\lambda,\mu}^{\lfloor \frac{i+1}{2} \rfloor}$, where $\lfloor \ell \rfloor$ is the integral part of a real number ℓ .

The finer filtration (1.3) on the space $\mathcal{D}_{\lambda,\mu}$ (clearly stable under the $\mathcal{K}(1)$ -action) is given by the spaces $\mathcal{D}_{\lambda,\mu}^{\frac{\ell}{2}}$ of differential operators (4.2).

4.2 Conjugation of differential operators

There exists a $\mathcal{K}(1)$ -invariant *conjugation* map $*$: $\mathcal{D}_{\lambda,\mu} \rightarrow \mathcal{D}_{\frac{1}{2}-\mu, \frac{1}{2}-\lambda}$ defined by

$$\langle A\phi, \psi \rangle = (-1)^{p(A)p(\phi)} \langle \phi, A^*\psi \rangle \quad (4.5)$$

for any $A \in \mathcal{D}_{\lambda,\mu}$ and $\phi \in \mathcal{F}_\lambda$, $\psi \in \mathcal{F}_{\frac{1}{2}-\mu}$, where $\langle \cdot, \cdot \rangle$ is the bilinear form (3.2).

Clearly, $*$ is a $\mathcal{K}(1)$ -isomorphism $\mathcal{D}_{\lambda,\mu}^\ell \cong \mathcal{D}_{\frac{1}{2}-\mu, \frac{1}{2}-\lambda}^\ell$ for every $\ell \in \frac{1}{2}\mathbb{N}$. In the particular case $\lambda + \mu = \frac{1}{2}$, the map $*$ is an involution. The module $\mathcal{D}_{\lambda, \frac{1}{2}-\lambda}^\ell$ splits into a direct sum of the submodules of *symmetric* and *skew-symmetric* operators.

The explicit formula of the conjugation map is easy to calculate:

$$* : \overline{D}^k \mapsto (-1)^{\lfloor \frac{k+1}{2} \rfloor} \overline{D}^k. \quad (4.6)$$

4.3 The principal symbol map

The highest order coefficient in (4.2) has the following geometric meaning.

Proposition 4.4. *For every $\ell \in \frac{1}{2}\mathbb{N}$, we have*

$$\mathcal{D}_{\lambda,\mu}^\ell / \mathcal{D}_{\lambda,\mu}^{\ell-\frac{1}{2}} \cong \mathcal{F}_{\mu-\lambda-\ell}. \quad (4.7)$$

Proof. For a differential operator $A \in \mathcal{D}_{\lambda,\mu}^\ell$ given by (4.2), one easily checks that the expression

$$\sigma_{\text{pr}}(A) := a_\ell(x, \xi) \alpha^{\mu-\lambda-\ell} \quad (4.8)$$

is a well-defined $(\mu - \lambda - \ell)$ -density. \square

The $\mathcal{K}(1)$ -invariant projection

$$\sigma_{\text{pr}} : \mathcal{D}_{\lambda,\mu}^\ell \rightarrow \mathcal{F}_{\mu-\lambda-\ell} \quad (4.9)$$

will be called the *principal symbol map*.

4.4 Space of symbols of differential operators

Consider the graded $\mathcal{K}(1)$ -module $\text{gr}\mathcal{D}_{\lambda,\mu}$ associated with the filtration (1.3). Proposition 4.4 implies that this $\mathcal{K}(1)$ -module is a direct sum of density modules:

$$\text{gr}\mathcal{D}_{\lambda,\mu} = \bigoplus_{i=0}^{\infty} \mathcal{F}_{\mu-\lambda-\frac{i}{2}}.$$

Note that this module depends only on the shift, $\mu - \lambda$, of the weights and not on μ and λ independently. We call this $\mathcal{K}(1)$ -module the *space of symbols of differential operators* and denote it $\mathcal{S}_{\mu-\lambda}$. The space of ℓ -th order symbols is

$$\mathcal{S}_{\mu-\lambda}^\ell := \bigoplus_{i=0}^{2\ell} \mathcal{F}_{\mu-\lambda-\frac{i}{2}}, \quad \text{where } \ell \in \frac{1}{2}\mathbb{N}. \quad (4.10)$$

5 Non-trivial cohomology classes of Lie superalgebra $\mathfrak{osp}(1|2)$

Given a Lie (super)algebra and its module, what is the corresponding cohomology ring? We give here some partial information about $H^0(\mathfrak{osp}(1|2); \mathcal{D}_{\lambda, \mu})$ (the space of invariants) and $H^1(\mathfrak{osp}(1|2); \mathcal{D}_{\lambda, \mu})$. More precisely, we exhibit non-trivial cocycles and conjecture that these cocycles generate the respective cohomology spaces.

Proposition 5.1. *For every $k = 1, 3, 5, \dots$, the differential operator*

$$\overline{D}^k : \mathcal{F}_{\frac{1-k}{4}} \rightarrow \mathcal{F}_{\frac{1+k}{4}} \quad (5.1)$$

is $\mathfrak{osp}(1|2)$ -invariant.

The operators (5.1) were found in [11, 12], as analogs of the Bol operators on S^1 , (see, e.g., [24]). One can prove that Proposition 5.1 provides the complete list of $\mathfrak{osp}(1|2)$ -invariant operators on weighted densities, but we will not need this.

One more special property of the modules $\mathcal{D}_{\lambda, \mu}$ with (λ, μ) given by (1.2) is as follows.

Theorem 5.2. *For every $k = 1, 3, 5, \dots$, the linear map $\gamma_k : \mathfrak{osp}(1|2) \rightarrow \mathcal{D}_{\frac{1-k}{4}, \frac{1+k}{4}}$, defined by*

$$\gamma_k(X_f) = D^3(f) \overline{D}^{k-1} + \frac{k-1}{2} D^4(f) \overline{D}^{k-2}, \quad (5.2)$$

where f is the contact Hamiltonian of an element $X_f \in \mathfrak{osp}(1|2)$, is a non-trivial 1-cocycle.

Proof. To prove that γ_k is a 1-cocycle, consider a map $\gamma_k^\lambda : \mathfrak{osp}(1|2) \rightarrow \mathcal{D}_{\lambda, \lambda + \frac{k}{2}}$ defined by the same formula (5.2) for arbitrary λ . One checks that

$$L_{X_f}^{\lambda + \frac{k}{2}} \circ \overline{D}^k - (-1)^{p(f)} \overline{D}^k \circ L_{X_f}^\lambda = \left(\lambda + \frac{k-1}{4}\right) \gamma_k(X_f).$$

Hence, $\gamma_k^\lambda = \frac{4}{4\lambda + k - 1} \delta \overline{D}^k$, so it satisfies the 1-cocycle condition if $\lambda \neq \frac{1-k}{4}$. By continuity, this is also true for $\lambda = \frac{1-k}{4}$.

The 1-cocycle γ_k^λ is a coboundary for any $\lambda \neq \frac{1-k}{4}$. Let us prove that in the case $\lambda = \frac{1-k}{4}$, this cocycle is indeed non-trivial.

We will need the explicit formula for γ_k in terms of the basis of $\mathfrak{osp}(1|2)$: it vanishes on all elements X_f of $\mathfrak{osp}(1|2)$ except $f = x^2$ or $x\xi$, for which

$$\begin{aligned} \gamma_k(X_{x^2}) &= 2\xi \overline{D}^{k-1} + \frac{k-1}{2} \overline{D}^{k-2}, \\ \gamma_k(X_{x\xi}) &= \overline{D}^{k-1}. \end{aligned} \quad (5.3)$$

Assume that there exists an operator $A \in \mathcal{D}_{\frac{1-k}{4}, \frac{1+k}{4}}$ such that γ_k is equal to δA , where

$$(\delta A)(X_f) = L_{X_f}^{\frac{1+k}{4}} \circ A - (-1)^{p(f)} A \circ L_{X_f}^{\frac{1-k}{4}}.$$

The operator A is of the form $A = a_m \overline{D}^m + \dots + a_0$; its principal symbol is a density of degree $\mu - \lambda - \frac{m}{2} = \frac{k-m}{2}$. If $m > k$, then, the principal symbol of $(\delta A)(X_f)$ is not identically zero since there is no $\mathfrak{osp}(1|2)$ -invariant density (except the constant function). Therefore, $m \leq k$ and $A = a_k \overline{D}^k + \dots + a_0$, where $a_k = \text{const}$. According to Proposition 5.1, the term \overline{D}^k commutes with the action of $\mathfrak{osp}(1|2)$, so it remains to consider $A = a_{k-1} \overline{D}^{k-1} + \dots + a_0$ with $a_{k-1} \neq 0$.

The principal symbol of A is a $\frac{1}{2}$ -density $a_{k-1} \alpha^{\frac{1}{2}}$. By assumption, the principal symbol of the operator is

$$\sigma_{\text{pr}}(\gamma_k(X_f)) = L_{X_f}^{\frac{1}{2}} \left(a_{k-1} \alpha^{\frac{1}{2}} \right)$$

Finally, the relation $\gamma_k(X_f) = 0$ implies that a_{k-1} is constant contradicting (5.3). \square

Remark 5.3. 1) The cocycle γ_k is odd (since k in (5.2) is odd).

2) One checks that

$$\gamma_k(X_f)^* = (-1)^{\frac{k-1}{2}} \gamma_k(X_f) \quad \text{for each } X_f \in \text{osp}(1|2).$$

Computation of the cohomology of $\text{osp}(1|2)$ with coefficients in $\mathcal{D}_{\lambda,\mu}$ is an interesting open problem. We formulate a conjecture on the structure of the first cohomology space.

Conjecture 5.4. *One has $H^1(\text{osp}(1|2); \mathcal{D}_{\lambda,\mu}) = \mathbb{C}^{0|1}$ if and only if $(\lambda, \mu) = (\frac{1-k}{4}, \frac{1+k}{4})$, with odd k and spanned by the cocycle γ_k . Otherwise, this cohomology space is trivial.*

In the case of S^1 and the Lie algebra $\mathfrak{sl}(2)$ a similar result was obtained in [17].

6 Equivariant quantization and symbol maps

We restrict the $\mathcal{K}(1)$ -action on $\mathcal{D}_{\lambda,\mu}^\ell$ to the subalgebra $\text{osp}(1|2)$ and look for an $\text{osp}(1|2)$ -isomorphism between $\mathcal{D}_{\lambda,\mu}^\ell$ and $\mathcal{S}_{\mu-\lambda}^\ell$ providing a “total symbol” of differential operators. We prove existence and uniqueness (up to normalization) of such an isomorphism for generic (λ, μ) and investigate the resonant case.

6.1 The main theorem

A map $\sigma : \mathcal{D}_{\lambda,\mu} \rightarrow \mathcal{S}_{\mu-\lambda}$ is called a *symbol map* if it is bijective and for every $\ell \in \frac{1}{2}\mathbb{N}$ the following diagram is commutative:

$$\begin{array}{ccc} \mathcal{D}_{\lambda,\mu}^\ell & \xrightarrow{\sigma} & \mathcal{S}_{\mu-\lambda}^\ell \\ \sigma_{\text{pr}} \downarrow & & \downarrow \\ \mathcal{F}_{\mu-\lambda-\ell} & \xrightarrow{\text{Id}} & \mathcal{F}_{\mu-\lambda-\ell} \end{array} \quad (6.1)$$

where the right arrow is the projection. In other words, the highest-order term of σ coincides with the principal symbol map. The inverse map, $Q = \sigma^{-1}$, is called the *quantization map*.

Recall that we call (λ, μ) *non-resonant* if $\mu - \lambda$ is not as in (1.1). The following statement is the main result of this paper; the proof will be given in the next section.

Theorem 6.1. *(i) If $\mu - \lambda$ is non-resonant, then $\mathcal{D}_{\lambda,\mu} \cong \mathcal{S}_{\mu-\lambda}$ as $\text{osp}(1|2)$ -modules, the isomorphism being given by the unique $\text{osp}(1|2)$ -invariant symbol map*

$$\sigma_{\lambda,\mu}(A) = \sum_{n=0}^k (-1)^{\lfloor \frac{n+1}{2} \rfloor} \frac{\binom{\lfloor \frac{k}{2} \rfloor}{\lfloor \frac{2n+1-(-1)^{n+k}}{4} \rfloor} \binom{\lfloor \frac{k-1}{2} \rfloor + 2\lambda}{\lfloor \frac{2n+1+(-1)^{n+k}}{4} \rfloor}}{\binom{2(\mu-\lambda) + n - k - 1}{\lfloor \frac{n+1}{2} \rfloor}} D^n(a) \alpha^{\mu-\lambda+\frac{n-k}{2}} \quad (6.2)$$

where $A = a(x, \xi) \overline{D}^k \in \mathcal{D}_{\lambda, \mu}^{\frac{k}{2}}$.

(ii) In the resonant case the $\text{osp}(1|2)$ -modules $\mathcal{D}_{\lambda, \mu}$ and $\mathcal{S}_{\mu - \lambda}$ are not isomorphic, except if (λ, μ) are given by (1.2).

Note that the binomial coefficients in (6.2) are defined by $\binom{\nu}{q} = \frac{\nu(\nu-1)\cdots(\nu-q+1)}{q!}$. This expression makes sense for arbitrary $\nu \in \mathbb{C}$.

Remark 6.2. The map $\sigma_{\lambda, \mu}$ is defined with the help of local coordinates (x, ξ) . Nevertheless, this map is independent with respect to the fractional-linear coordinate transformations (2.10). In particular, if one fixes a projective structure on $S^{1|1}$, the map $\sigma_{\lambda, \mu}$ is globally defined. This follows from the $\text{osp}(1|2)$ -equivariance.

Let us now give the explicit formula for the quantization map.

Proposition 6.3. The map $Q_{\lambda, \mu} = \sigma_{\lambda, \mu}^{-1}$ associates to a tensor density $\varphi = f \alpha^{\mu - \lambda - \frac{k}{2}}$ the following differential operator from \mathcal{F}_λ to \mathcal{F}_μ :

$$Q_{\lambda, \mu}(\varphi) = \sum_{n=0}^k \frac{\binom{\left\lfloor \frac{k}{2} \right\rfloor}{\left\lfloor \frac{2n+1-(-1)^{n+k}}{4} \right\rfloor} \binom{\left\lfloor \frac{k-1}{2} \right\rfloor + 2\lambda}{\left\lfloor \frac{2n+1+(-1)^{n+k}}{4} \right\rfloor}}{\binom{2(\mu - \lambda) - k + \left\lfloor \frac{n-1}{2} \right\rfloor}{\left\lfloor \frac{n+1}{2} \right\rfloor}} D^n(f) \overline{D}^{k-n}. \quad (6.3)$$

Remark 6.4. In [4], the $\text{osp}(1|2)$ -quantization map was written down in the particular case $\lambda = \mu = 0$. The authors use the form (4.3) of a differential operator. One can rewrite formula (6.3) using the same notations:

$$Q_{\lambda, \mu}(f \alpha^{\mu - \lambda - \frac{k}{2}}) = \sum_{n \geq 0} \frac{\binom{\left\lfloor \frac{k}{2} \right\rfloor}{\left\lfloor \frac{n+1}{2} \right\rfloor} \binom{\left\lfloor \frac{k-1}{2} \right\rfloor + 2\lambda}{\left\lfloor \frac{n}{2} \right\rfloor}}{\binom{k + 2(\lambda - \mu)}{\left\lfloor \frac{n+1}{2} \right\rfloor}} D^n(f) D^{k-n},$$

where the upper bound for summation is $2 \left\lfloor \frac{k+1}{2} \right\rfloor$. In the particular case $\lambda = \mu = 0$, this formula coincides with formula (7.5) of [4].

6.2 Proof of Theorem 6.1

The proof of Theorem 6.1 consists of three parts. First, we show that the symbol map (6.2) is, indeed, $\text{osp}(1|2)$ -equivariant. We also prove the existence for the special values (1.2). Then we prove the uniqueness of the symbol map (6.2). Finally, we show that, for the resonant values of $\mu - \lambda$, there is no isomorphism between $\mathcal{D}_{\lambda, \mu}$ and $\mathcal{S}_{\mu - \lambda}$ if λ and μ are not given by (1.2).

A. Existence and explicit formula. Consider symbol maps given by differential operators of the form

$$\sigma(A) = \sum_{n=0}^k \beta_n^k(x, \xi) D^n(a) \alpha^{\mu - \lambda + \frac{n-k}{2}}, \quad (6.4)$$

where $A = a(x, \xi) \overline{D}^k$, and $\beta_n^k(x, \xi) \in C_c^\infty(S^{1|1})$ are arbitrary functions, and calculate the condition of $\text{osp}(1|2)$ -equivariance. Clearly, it suffices to impose invariance with respect to D and $x D$ and the following is straightforward.

Lemma 6.5. (i) A symbol map (6.4) commutes with the action of D if and only if the coefficients $\beta_n^k(x, \xi)$ are constants (i.e., do not depend on x, ξ and on the parity of A).

(ii) A symbol map (6.4) commutes with the action of xD if and only if the following system is satisfied:

$$\begin{aligned} p\beta_{2m-1}^{2s-1} &= m\beta_{2m}^{2s}, \\ s\beta_{2m}^{2s-1} &= (2(\lambda - \mu) - m + 2s)\beta_{2m+1}^{2s}, \\ (2\lambda + s)\beta_{2m-1}^{2s} &= m\beta_{2m}^{2s+1}, \\ (2\lambda + s)\beta_{2m}^{2s} &= (2(\lambda - \mu) - m + 2s + 1)\beta_{2m+1}^{2s+1}, \end{aligned} \tag{6.5}$$

where $1 \leq s \leq \lfloor \frac{k}{2} \rfloor$ and $0 \leq m \leq s$.

If $\mu - \lambda$ is non-resonant, then it is easy to see that the solution of the system (6.5) with the initial condition $\beta_0^k = 1$ (which is equivalent to the fact that σ preserves the principal symbol) is unique and given by (6.2).

If $\mu - \lambda$ is resonant and given by (1.2), then the system (6.5) can also be easily solved. The solution is no longer unique since the system is split into separate independent parts. Note that in this case, the explicit solution can still be obtained from formula (6.2) if one chooses an arbitrary resolution of 0/0-singularities.

B. Uniqueness. Consider the non-resonant case. The main ingredient of our proof of the uniqueness of an $\mathfrak{osp}(1|2)$ -equivariant symbol map is the *locality* property and Definition 1 of differential operators, see Section 4.1. This idea is borrowed from [18]. Since any $\mathfrak{sl}(2)$ -equivariant map

$$T : \mathcal{F}_{\nu_1}(S^1) \rightarrow \mathcal{F}_{\nu_2}(S^1), \tag{6.6}$$

with $\nu_1 \geq \nu_2$, is local (see Theorem 5.1 from [18]), so is the symbol map. Indeed, given two $\mathfrak{osp}(1|2)$ -equivariant symbol maps σ_1 and σ_2 , the map

$$\sigma_2^{-1} \circ \sigma_1 : \mathcal{S}_{\mu-\lambda} \rightarrow \mathcal{S}_{\mu-\lambda}$$

is $\mathfrak{osp}(1|2)$ -equivariant. Decomposition (4.10) shows that $\sigma_2^{-1} \circ \sigma_1$ is a sum of maps (6.6). Therefore, $\sigma_2^{-1} \circ \sigma_1$ is a differential operator. It follows that the general form of a symbol map σ commuting with the action of $\mathfrak{osp}(1|2)$ is given by formula (6.4). However, we already proved in part **A** that such a symbol map is unique and given by (6.2).

C. Cohomological obstructions. Assume now that the shift of the weight is resonant, $\mu - \lambda = \frac{m}{2}$, but (λ, μ) are not as in (1.2). We will prove that there is no $\mathfrak{osp}(1|2)$ -isomorphism between $\mathcal{D}_{\lambda, \mu}$ and $\mathcal{S}_{\mu-\lambda}$ in this case. Indeed, recall ([7], Section 1.4.5) that any exact sequence of \mathfrak{g} -modules

$$0 \longrightarrow V \xrightarrow{i} W \longrightarrow U \longrightarrow 0,$$

defines an element in $H^1(\mathfrak{g}; \text{Hom}(U, V))$: for any section $\tau : U \rightarrow W$, we define a 1-cocycle on \mathfrak{g} by

$$c_\tau(X)(u) = i^{-1}((\rho_X \circ \tau - (-1)^{p(\tau)p(\rho_X)} \tau \circ \rho_X)(u)),$$

where $u \in U$ and ρ_X is the action of $X \in \mathfrak{g}$; the cohomology class $[c_\tau]$ of c_τ is independent of the choice of τ . The sequence is split if and only if $[c_\tau] = 0$.

The exact sequence of $\mathfrak{osp}(1|2)$ -modules

$$0 \longrightarrow \mathcal{D}_{\lambda, \mu}^{\frac{m-1}{2}} \xrightarrow{i} \mathcal{D}_{\lambda, \mu}^{\frac{m}{2}} \xrightarrow{\sigma_{\text{pr}}} \mathcal{F}_0 \longrightarrow 0, \tag{6.7}$$

where i is the inclusion, defines an element of $H^1(\mathfrak{osp}(1|2); \text{Hom}(\mathcal{F}_0, \mathcal{D}_{\lambda, \mu}^{\frac{m-1}{2}}))$; the principal symbol map $\sigma_{\text{pr}} : \mathcal{D}_{\lambda, \mu}^{\frac{m-1}{2}} \rightarrow \mathcal{F}_{\frac{1}{2}}$ defines $\bar{c} \in H^1(\mathfrak{osp}(1|2); \text{Hom}(\mathcal{F}_0, \mathcal{F}_{\frac{1}{2}}))$. Set: $\tau : a \mapsto a \bar{D}^m$. The following lemma is a straightforward computation.

Lemma 6.6. *If m is odd, then $\bar{c} = (\lambda + \frac{m-1}{4}) \gamma_1$; if m is even, then $\bar{c} = \frac{m}{2} \gamma_1$.*

If $\bar{c} \neq 0$, then the sequence (6.7) is not split. The module $\mathcal{D}_{\lambda, \mu}^{\frac{m}{2}}$ is then non-isomorphic to the corresponding graded module of symbols. Theorem 6.1 is proved. \square

Proof of Proposition 6.3 is similar to that of Theorem 6.1, part (i).

6.3 Discussion

It is quite clear that the $\mathcal{K}(1)$ -modules $\mathcal{D}_{\lambda, \mu}^\ell$ and $\mathcal{S}_{\mu-\lambda}^\ell$ are not isomorphic for $\ell > 2$ and there is no $\mathcal{K}(1)$ -invariant symbol map. For instance, one can check that the unique $\mathfrak{osp}(1|2)$ -invariant symbol map from Theorem 6.1 does not commute with any other element of $\mathcal{K}(1)$, implying the above statement in the non-resonant case. However, we do not give here a complete proof. Note that a similar result holds in the case of S^1 , see [8, 10] and the multi-dimensional case [18]. This problem is related to the first cohomology space $H^1(\mathcal{K}(1); \mathcal{D}_{\lambda, \mu}^\ell)$ (cf. [1] for preliminary results on this subject).

We believe that the $\mathcal{K}(1)$ -modules $\mathcal{D}_{\lambda, \mu}^\ell$ deserve further study. We formulate here the problem of classification of these modules (see [8, 10] for the case of S^1), as well as existence of the exceptional weights (λ, μ) in the sense of [5]. It would also be interesting to study the corresponding automorphism groups (see [9]).

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